

# Binary black hole shadows, chaotic scattering and the Cantor set

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## Motivation

Using the **Event Horizon Telescope**, astronomers hope to generate high-resolution images of the black hole (BH) **shadow** at the centre of our galaxy in 2017.

Binary BH shadows can be understood by looking at **scattering of light** (null geodesics) around a pair of BHs.

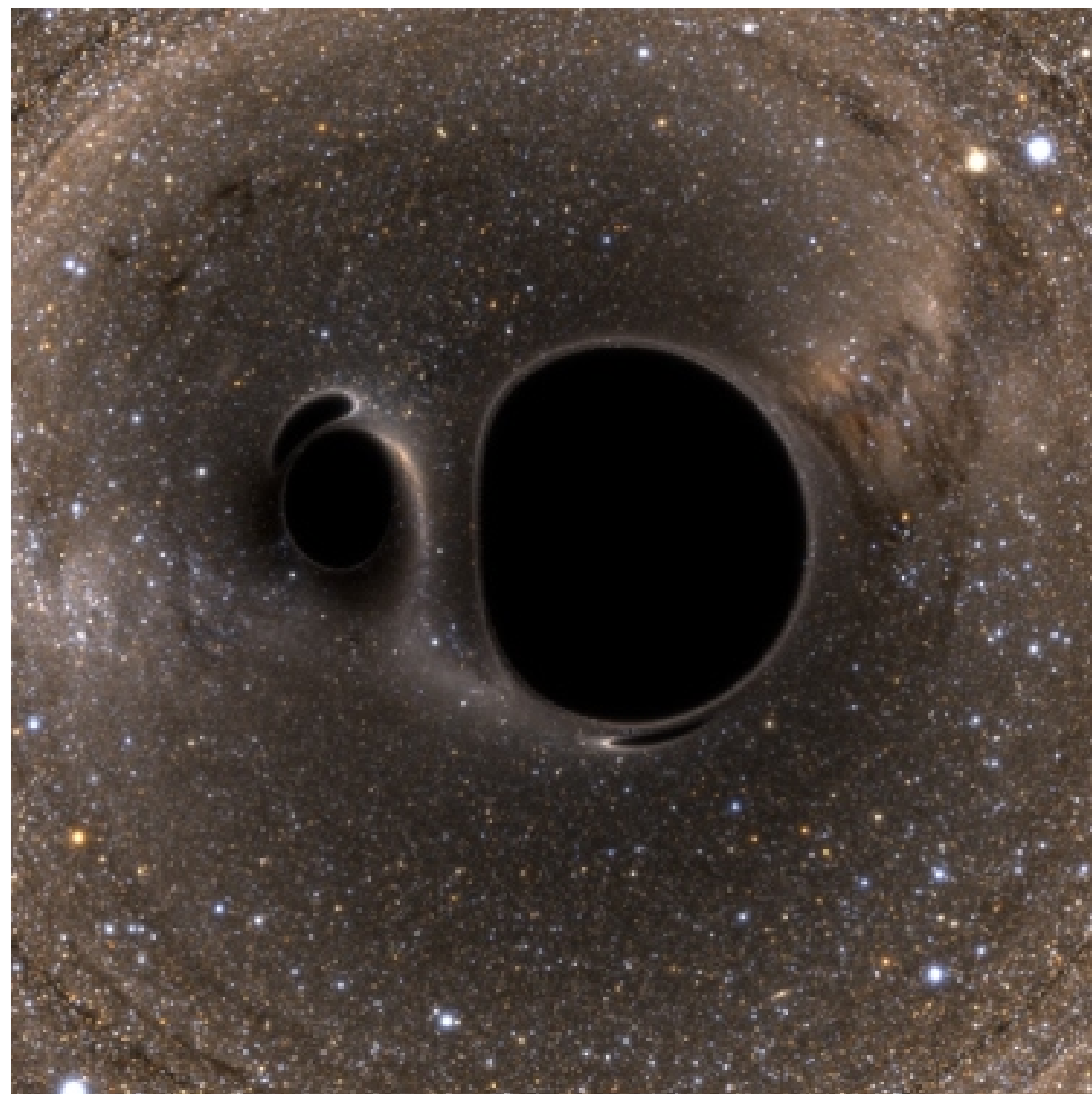


Figure 1: Binary BH shadow with self-similar eyebrow-like features. [Credit: SXS Lensing Group, www.black-holes.org]

## Model

We use a **toy model** of two extremal BHs in static equilibrium. The Majumdar–Papapetrou metric reads

$$ds^2 = -U^{-2}dt^2 + U^2dx \cdot dx, \quad (1)$$

where

$$U = 1 + \frac{M_1}{\sqrt{\rho^2 + (z - a/2)^2}} + \frac{M_2}{\sqrt{\rho^2 + (z + a/2)^2}}, \quad (2)$$

and  $a$  is the separation between the BHs.

The Hamiltonian for this system is

$$H = \frac{1}{2}\tilde{g}^{ab}p_ap_b = \frac{1}{2}(p_\rho^2 + p_z^2) + \frac{1}{2}\left(\frac{p_\phi^2}{\rho^2} - U^4\right). \quad (3)$$

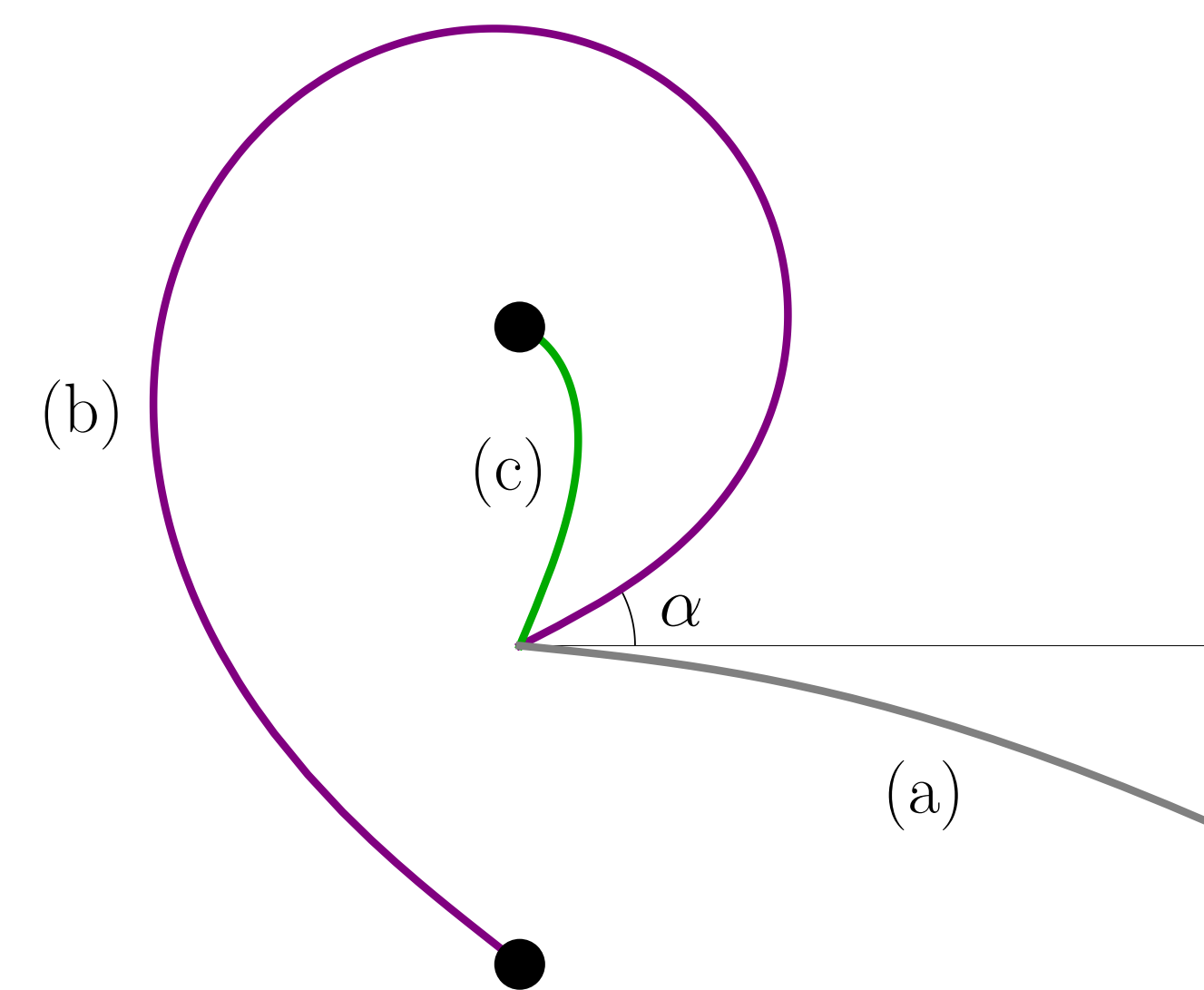
This is a **non-integrable 2D Hamiltonian system**. To analyse the scattering of light, we look at Hamilton's equations.

## What is a black hole shadow?

**Ray-casting approach:** Null rays are traced away from the camera lens, backwards in time. The shadow is the region of the image taken by the camera where rays are traced back to a BH.

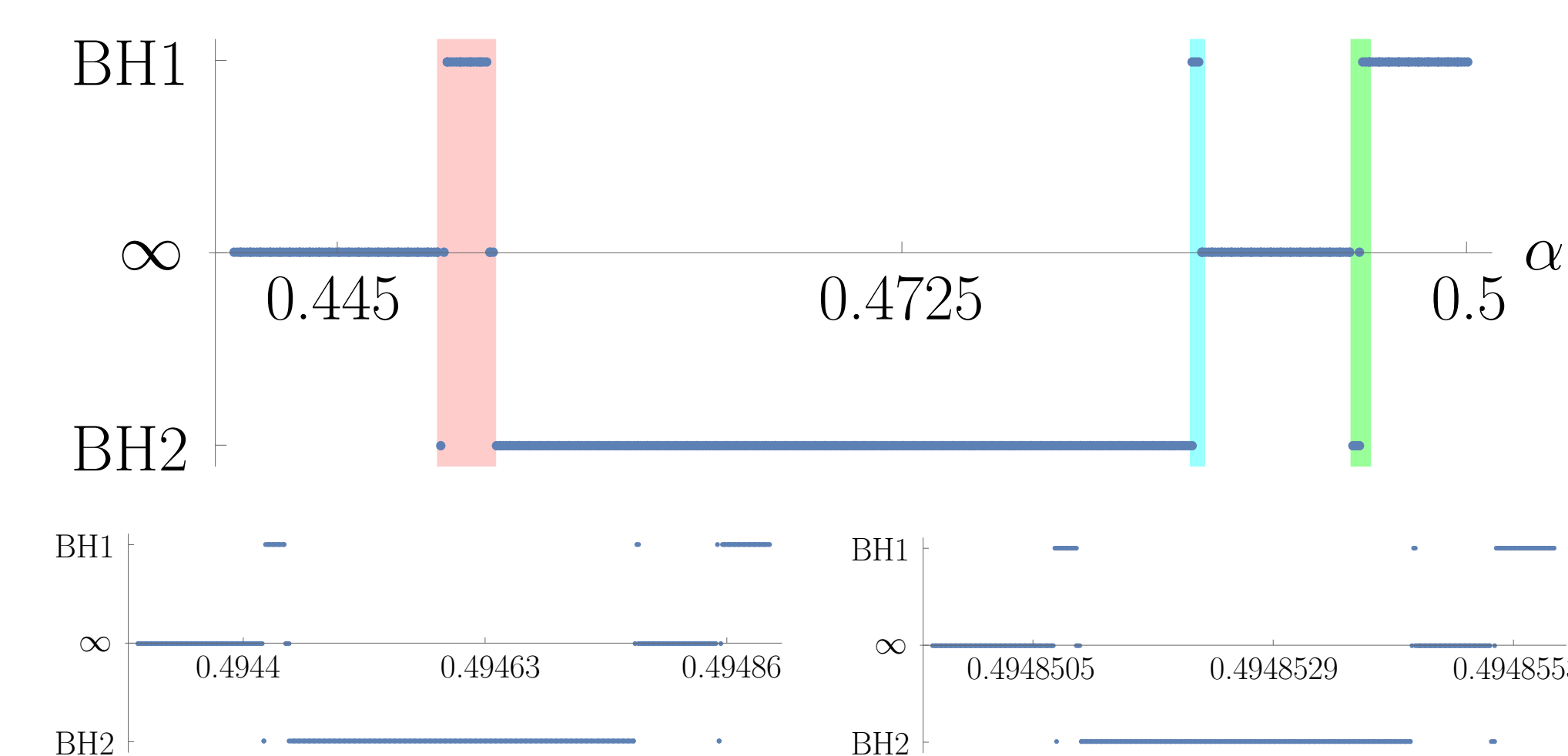
## One-parameter scattering problem

For simplicity, we look at a 1-parameter scattering problem. Send null rays out from midpoint between BHs, varying initial angle of elevation from horizontal. Null geodesics can: **(a)** escape to  $\infty$ ; **(b)** fall into lower BH; **(c)** fall into upper BH.



There is also a set of **unstable perpetual orbits** (not scattered or absorbed by the BHs). These lie on the interface between the absorbed and scattered rays.

The 1-parameter scattering problem admits a **1D shadow**. The structure of this 1D shadow allows us to understand 2D shadow images.

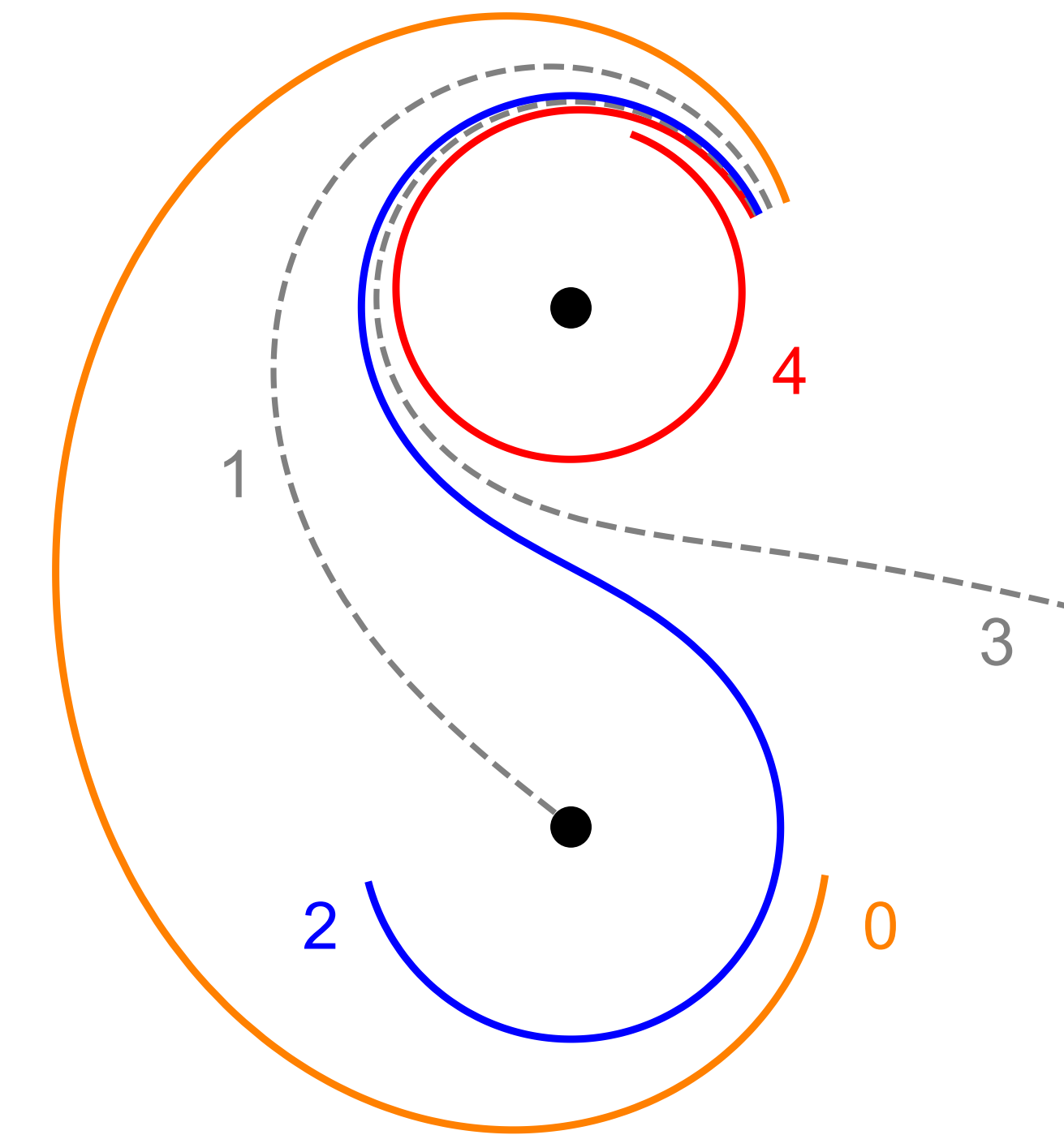


Zooming in on the region highlighted in green, we see **self-similarity** in the 1D shadow.

## Symbolic dynamics

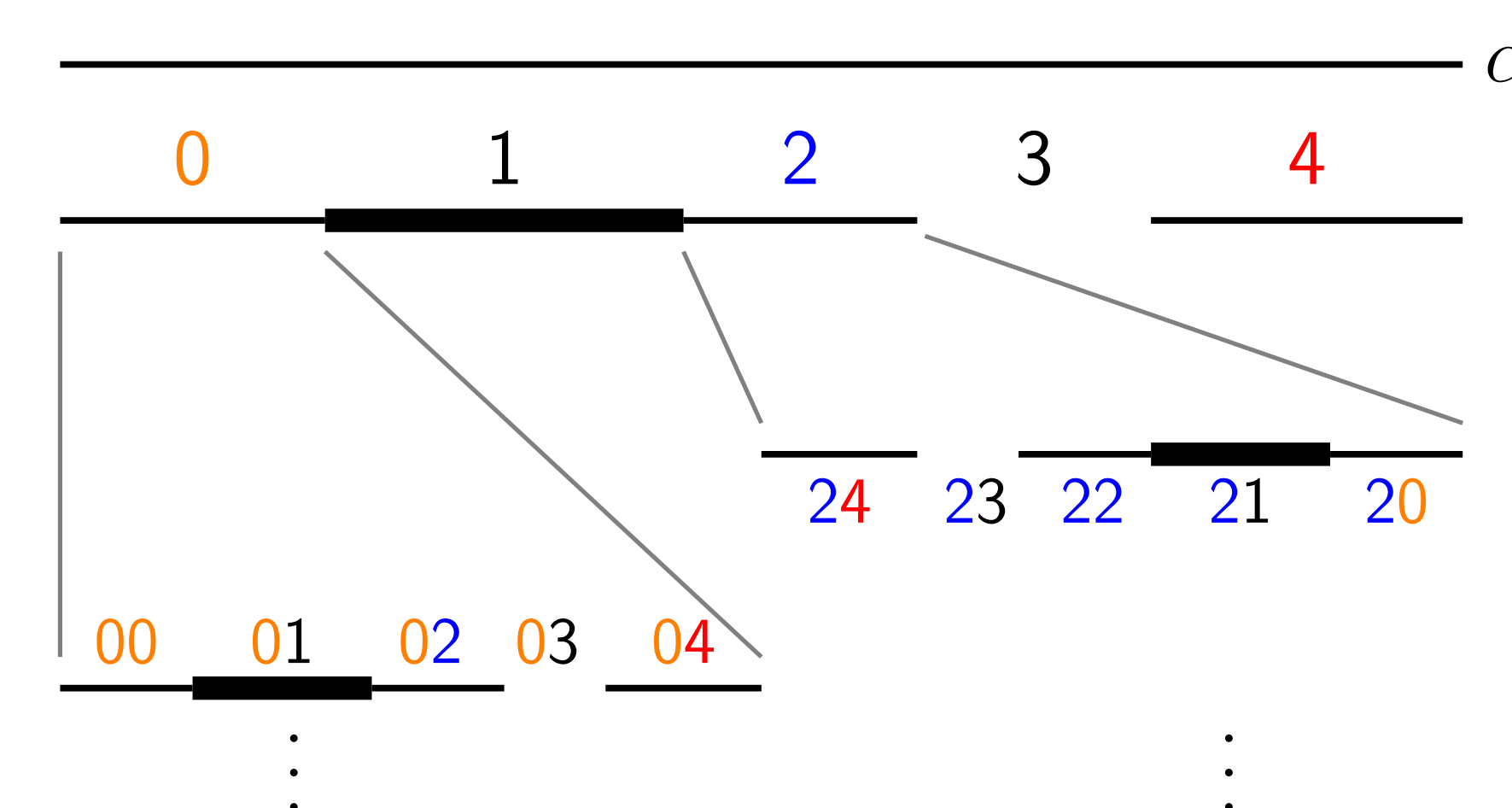
We have developed a **symbolic code** based on '**decision points**'. A geodesic which has travelled around a BH can:

- 0 travel around the other BH in the same sense;
- 1 fall into a BH (absorbed);
- 2 travel around the other BH in the opposite sense;
- 3 escape to  $\infty$  (scattered);
- 4 travel around the same BH again.



## Building the 1D shadow

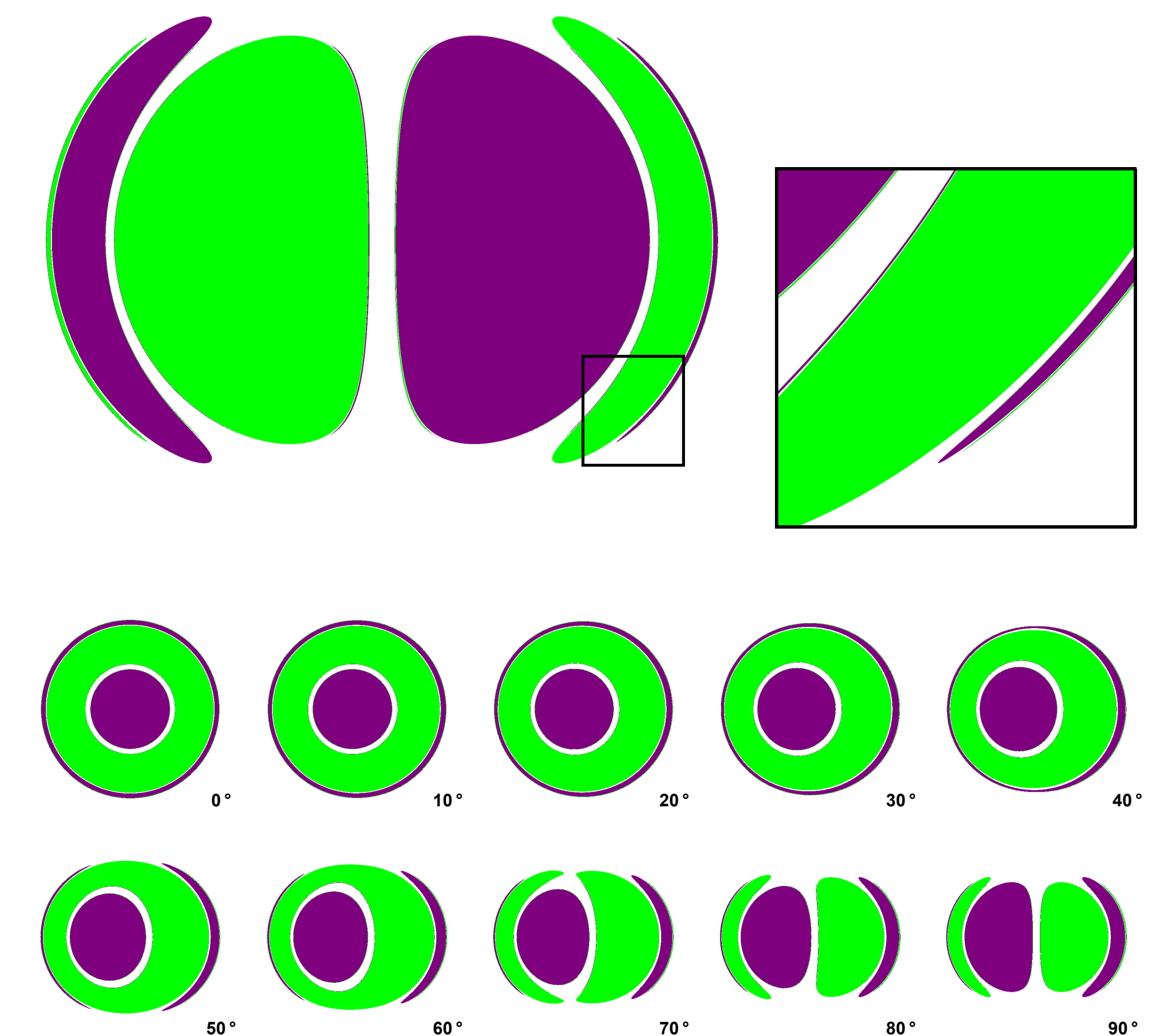
Using an iterative process, based on symbolic dynamics, we can construct 1D shadow. At each stage, remove intervals of initial data corresponding to decisions **1** and **3**. Iterate on the remaining intervals.



After infinitely many iterations, we are left with a **Cantor-like set** of perpetual orbits. The 1D shadow is the infinite union of open intervals, ending in the digit **1**.

## 2D shadows

The fractal structure of the 1D shadow is inherited by 2D shadows. We present a gallery of 2D shadows for different viewing angles. At  $0^\circ$  we look along axis connecting BHs. At  $90^\circ$  we look side-on at BHs.



In each case, main shadow images are surrounded by an infinite hierarchy of self-similar concentric ring-like structures or eyebrow-like features.

## Conclusions

- Binary BH system exhibits **chaotic scattering**: it admits an infinite set of perpetual orbits as a **Cantor-like set** on the initial data.
- We have developed **symbolic dynamics** to describe **perpetual orbits**.
- We were able to construct the BH shadow using an iterative process based on **symbolic dynamics**.
- Binary BH shadows have **fractal** properties.

## References

- [1] J. O. Shipley & S. R. Dolan [arXiv:1603.04469](https://arxiv.org/abs/1603.04469)

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